Impact of Atmospheric Turbulence on SLR system

AZU

NESC Meeting

Julien Chabé

With contributions from my colleagues : Duy Hà Phung, Clément Courde, Aziz Ziad, Eric Aristidi.

Université Côte d'Azur, Observatoire de la Côte d'Azur, CNRS, IRD, Géoazur, 2130 Route de l'Observatoire, 06460 Caussols, France;

1

What I will talk about

- What is atmospheric turbulence and how to measure it ?
- What is the effect for SLR systems ?
- Upcoming challenges regarding turbulence
- Conclusion



How to quantify turbulence ?

- Atmosphere is stratified in turbulent layers which are in constant evolution...
- A local parameter to quantify the <u>turbulence</u> « strength » : C_n^2 in m^{-2/3}
- By computing the refractive index structure function:

n(r)

Refraction index at 2 positions :

• one can found:

77U

$$D_n(\vec{\rho}) = \left\langle \left[\delta n(\vec{r}) - \delta n(\vec{r} + \vec{\rho})\right]^2 \right\rangle = C_n^2 \rho^{2/3}$$

• C_n^2 is the Refraction index structure constant (If no turbulence : $\delta n(r)$ - $\delta n(r+\rho)$ constant : $C_n^2=0$



Grasse Observatory (Moon Limb profiler)

How to quantify turbulence ?

And the Fried Parameter (r_0) and astronomical Seeing (ϵ)

$$r_0 = \left[0.42k^2 \int_0^L C_n^2(z) \, dz\right]^{-3/5}$$



ε = 0"3

ε = 1"

ε = 0"1

$$\epsilon = 0.98 \frac{\lambda}{r_0}$$

• The isoplanetic angle θ_0 is the angular coherence of the turbulence:

$$\theta_0 = \left[2.91k^2 \int_0^L C_n^2(z) z^{5/3} dz\right]^{-3/5}$$

- Typical values for θ_0 at wavelength λ =500nm are roughly **7–10 µrad** for a near vertical path

ε = 0"05

• The coherence time related to wind speed (V) :

$$\tau_0 = \left[2.91k^2 \int_0^L C_n^2(z) V^{5/3}(z) \, dz \right]^{-3/5} \qquad \tau_0 = \frac{0.32r_0}{V_\perp}.$$

How to quantify turbulence ?

 In Kolmogorov theory of turbulence : injection of energy occurs at a large scale : the outer scale L₀

- The outer scale L₀ is not easy to evaluate but it is important parameter often neglected...
- The outer scale L₀ it is reponsible for tip/tilt effect -> 80% of the turbulence



Shack-Hartmann analysis



- Downlink laser beam wavefront distortion analysis: The Shack-Hartmann Wave Front Sensor measures the local slope of incoming wavefront
- Fried parameter r₀ determination (radius of coherence of the wavefront)

Generalized Differential Image Monitor







Differential image motion of stars

- Independant atmospheric turbulence analysis:
 - Fried parameter r₀
 - Scintillation index
 - Isoplanatic angle
 - Outer scale of turbulence: L₀

[Aristidi 2019] E Aristidi, A Ziad, J Chabé, Y Fantéi-Caujolle, C Renaud, C Giordano, A generalized differential image motion monitor, Monthly Notices of the Royal Astronomical Society, Volume 486, Issue 1, June 2019, Pages 915–925, <u>https://doi.org/10.1093/mnras/stz854</u>

[Ziad 2016] Aziz Ziad. Review of the outer scale of the atmospheric turbulence. Adaptive Optics Systems V. SPIE, July 2016

C_n² (h) Profiler (Moon/Sun)



ATS





Principle : mesure the agitation at 2 points of the limb \rightarrow 2 points far away : access to low altitude turbulent layers \rightarrow 2 close points : high altitude layers

- Estimation of the turbulence profile C_n^2 (h) up to h=24km : 33 layers
- 1 profile every minute
- Resolution: 100m (low altitude) to 2km (high altitude)



Laser + atmosphere

 The coherence of the laser beam wavefront is destroyed by the turbulent atmosphere:



5

Laser + atmosphere

 The speckle pattern has coherence time of ~ms

77

- The instantenous center of the beam (hot spot) where the intensity is the strongest is moving randomly around the average central position
- When we average on a long time, the beam looks gaussian but is larger than the original one
- Scintillation of the light in the detector plane



Figure 6.7 (*a*) Beam wander as described by movement of the "hot spot" (instantaneous center) within the beam. (*b*) The long-term spot size is the result of beam wander, beam breathing, and diffraction. The shaded circles depict random motion of the short-term beam in the receiver plane.

Scintillation of the laser beam

Example of the scintillation of a telecom downlink laser beam

$$\sigma_I^2 = \frac{\sigma^2}{I^2} \qquad \sigma_I^2 = 17D^{-7/3}\cos(\theta_{zen})^{-3} \int_0^L h^2 C_n^2(h) dh$$





- The turbulence effects (scintillation, wavefront distortion) induce fades in the detection.
- These fades can occur 10s to 100s of times per second for vertical links between the ground and space
- This problem increases with small size detectors
- Good detectors with low timing jitter:
 - Small active area : tens of μm
 - Single mode fiber coupling (SMF 28 -> 10µm core diameter) for superconducting nanowire single-photon detector : D~10 µm at Delft university



ACS Photonics 2020, 7, 7, 1780–1787 https://doi.org/10.1021/acsphotonics.0c00433

Example for telecom downlink laser at Grasse

- Received Laser is splitted into:
 - 90% Laser com analysis:

- APD telecom : 80µm diameter (30 arcsec FOV)
- 10% for Fine tracking by TipTilt mirror
 - Spot laser position on TipTilt camera was stabilized with RMS < 1.1 arcsec during satellite pass, at 50 Hz from low (10 deg) to high (48 deg) elevation



Turbulence induced Ranging jitter for SLR

$$\left< \Delta L^2 \right> = 3.127 L_0^{(5/3)} \int C_n^2(h) dh$$

- Turbulence gives a piston like effect -> Timing jitter
- This effect is dependent on L₀ the outer scale of the turbulence (~10 to 100m)
- Few µm to mm (strong turbulence like horizontal path over several kms)



[Gardner 1976] C. S. Gardner. Effects of random path fluctuations on the accuracy of laser ranging systems. Appl. Opt., 15(10):2539–2545, Oct 1976 [Kral 2005] Lukas Kral, Ivan Prochazka, and Karel Hamal. Optical signal path delay fluctuations caused by atmospheric turbulence. Opt. Lett., 30(14):1767–1769, Jul 2005.

Turbulence induced Ranging jitter for SLR

 Continuous telecom laser (10Gbits) phase measurement at Grasse (ground target):



Coherent optical Doppler Orbitography & frequency transfer

- Nice demonstrations on ground target of phase measurement with active phase noise compensation + tip tilt mirror
 - Dix-Matthews, B.P., Schediwy, S.W., Gozzard, D.R. et al. Point-to-point stabilized optical frequency transfer with active optics. Nat Commun 12, 515 (2021). <u>https://doi.org/10.1038/s41467-020-20591-5</u>
 - Dix-Matthews, B.P., Schediwy, S.W., Gozzard, D.R. et al. Methods for coherent optical Doppler orbitography. J Geod 94, 55 (2020). https://doi.org/10.1007
- Ground to space experiment still to be demonstrated
 - Point ahead angle vs isoplanatic angle ?
 - Fades ?



Fig. 1 An overview of the technique that is being considered. An optical frequency (ν_L) is reflected off a retro-reflector located on the satellite, where it undergoes a Doppler shift (ν_D) dependent on the satellite's radial velocity (\overline{v}) . The satellite's radial velocity is then determined using the Doppler shifted return frequency. The transmitted and reflected signals also experience frequency perturbations caused by the atmospheric phase noise $(\Delta \phi_{atm})$.







Conclusion

- Atmospheric turbulence induces : Beam wandering, Beam spreading, Scintillation, fading, phase noise
- Instrumentation to measure turbulence exist:
 - Shack-Hartmann analyser: In situ measurement, easy to built/analyse, gives r₀
 - Differential Image Monitor: independant instrument (stars), easy to built/analyse, cheap, gives r_0 , L_0 , θ_0
 - C_n^2 (h) Profiler : Very High resolution of all the params.
- For advanced SLR (100kHz MHz, SNSPD):
 - Problematic with high performances (small) detectors and single mode fibers
 - A tip-tilt correction is a first step to mitigate this (~80% of the turbulence) and probably good enough for most « upgraded SLR system »
- Using the coherence properties of a laser beam: difficult but a lot of progress has been done
 - Some adaptive optics has be used to corrected wavefront degradation for single mode fiber injection
 - Active phase correction for coherent optical Doppler Orbitography & frequency transfer



GDIMM for ~5k€